

LATTICE-BOLTZMANN SCHEME FOR DENDRITIC GROWTH IN PRESENCE OF CONVECTION

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SUMMARY

A combined phase-field/lattice-Boltzmann scheme is proposed to simulate dendritic growth from supercooled melt, with account for flows of liquid and thermal convection.

MODEL

- 1. Simulation of solidification phase-field model (Karma and Rappel, 1996).
- 2. Flow of liquid lattice-Boltzmann-BGK (LBGK) method with incorporated interactions with solid and thermal convection. This step can be left out in the case of purely diffusional growth.
- 3. Conductive and convective heat transfer multicomponent LBGK method.

Phase-field model

$$\begin{split} \tau(\theta)\phi_t &= (2\phi-1-4\lambda\overline{T}\,\phi(1-\phi))2\phi(1-\phi) + \\ \nabla\cdot\left(W^2(\theta)\nabla\phi\right) - \partial_x\left(W(\theta)W'(\theta)\phi_y\right)) + \\ \partial_y\left(W(\theta)W'(\theta)\phi_x\right), \\ \overline{T}_t + \mathbf{U}\nabla\overline{T} &= D\nabla^2\overline{T} + \phi_t. \end{split}$$

$$\begin{split} \phi &- \text{concentration of solid phase } (0 \leq \phi \leq 1), \\ \overline{T} &= \rho c_p (T - T_m) / L - \text{normalized temperature,} \\ W &- \text{anisotropic interface width,} \\ \tau &- \text{relaxation time.} \end{split}$$

In order to obtain zero kinetic coefficient, the following relations must be imposed

$$W = W_0 A(\theta), \ \tau = \tau_0 A^2(\theta), \ \lambda = \frac{2ID\tau_0}{(K + JF)W_0^2}.$$

Anisotropy function

$$A(\theta) = 1 + \varepsilon \cos 4\theta,$$

 $\theta = \arctan(\phi_y/\phi_x)$ — the angle between the local interface normal and the X axis. We assume $\tau_0 = 1$, $W_0 = 1$. Interface stiffness $\alpha = 15\varepsilon$.

This equation was discretized on a uniform spatial lattice with a step $\Delta x = 0.4$, and solved using the explicit Euler method with time step Δt .

Lattice-Boltzmann method

- Regular lattice, lattice vectors \mathbf{e}_k
- Discrete set of velocities $\mathbf{c}_k \Delta t_{LB} = \mathbf{e}_k$

Variables f_k — one-particle distribution functions. Evolution equation

$$f_k(t, \mathbf{x}) = f_k(t - \Delta t_{LB}, \mathbf{x} - \mathbf{c}_k \Delta t_{LB}) + \frac{f_k^{eq} - f_k}{\tau_f}.$$

Propagation and collisions, τ_f — Maxwellian relaxation time. Later on $\Delta t_{LB} = 1$. Kinematic viscosity $\nu = (\tau_f - 1/2)/3$. Hydrodynamic quantities

$$\rho = \sum_{k} f_k, \qquad \rho \mathbf{U} = \sum_{k} f_k \mathbf{c}_k.$$

Equilibrium distribution functions $f_k^{eq}(\rho, \mathbf{U})$. Interaction with solid

$$\mathbf{F}_d = -\nu \frac{2h\phi^2}{W_0^2} \mathbf{U}.$$

Thermal convection — buoyancy force

$$\mathbf{F}_c = -\rho\alpha(1-\phi)(\overline{T}-\overline{T}_0)\mathbf{g},$$

 α — coefficient of thermal expansion, ${\bf g}$ — gravity acceleration.

Heat transport

Second set of distribution functions N_k . Evolution equation

$$N_{k}(t + \Delta t, \mathbf{x} + \mathbf{c}_{k}\Delta t) = N_{k}(t, \mathbf{x}) + \frac{N_{k}^{eq} - N_{k}}{\tau_{T}}.$$

$$N_{k}^{eq} = N_{k}^{eq}(\overline{T}, \mathbf{U} + \Delta \mathbf{U}/2)$$

$$\overline{T} = \sum_{k} N_{k}, \quad \mathbf{U} = \sum_{k} f_{k}\mathbf{c}_{k} / \sum_{k} f_{k}, \quad \Delta \mathbf{U} = \mathbf{F} / \sum_{k} f_{k}$$

$$\mathbf{F} = \mathbf{F}_{d} + \mathbf{F}_{c}$$
Total force.
Resulting equation

$$\frac{\partial \overline{T}}{\partial t} = \mathbf{U}\nabla \overline{T} + \chi \nabla^2 \overline{T} + \frac{\partial \phi}{\partial t}$$

Thermal diffusivity $\chi = (\tau_T - 1/2)/3$ everywhere (symmetric model).

Shear flow



Figure 1: Dendrite. $\Delta = 0.7, 15\varepsilon = 0.15$. Reduced velocity $\overline{U}=0$ (a), $\overline{U}=0.0123$ (b), $\overline{U}=0.0247$ (c), $\overline{U}=0.0493$ (d). Interface contours are shown at time increments of 50



Figure 2: Seaweed. $\Delta = 0.8, 15\varepsilon = 0.15$. Reduced velocity $\overline{U}=0$ (a), $\overline{U}=0.0247$ (b), $\overline{U}=0.0493$ (c), $\overline{U}=0.0987$ (d). Interface contours are shown at time increments of 20

Thermal convection



Figure 3: Dendritic growth with convection. $\Delta = 0.8, 15\varepsilon = 0.3$. $\beta g_y = 0.0, g_x = 0$ (a); $\beta g_y = -0.0005, g_x = 0$ (b); $\beta g_y = 0.0005, g_x = 0$ (c) and $g_y = 0, \beta g_x = 0.0005$ (d). Interface contours are shown at time increments of 10

Influence of parallel flow on the growth

Growth of the dendrite tip in the parallel flow

 $\Delta = 0.65, 15\varepsilon = 0.75, \text{ grid size } 300 \times 600$



Figure 4: Dependence of reduced velocity \overline{V} and selection parameter σ on flow Reynolds number. $1 - \overline{V}, \nu = 1/3; 2 - \sigma, \nu = 1/3; 3 - \overline{V}, \nu = 1/6; 4 - \sigma, \nu = 1/6$



Figure 5: Dependence of reduced tip radius $\overline{\rho}$ on reduced flow velocity \overline{U}



Figure 6: Dependence of reduced velocity \overline{V} and selection parameter σ on reduced flow velocity \overline{U}



Figure 7: Dependence of reduced tip radius $\overline{\rho}$ on reduced flow velocity \overline{U}

Dependences $\overline{V} \sim \overline{U}^{0.38}$, $\overline{R} \sim \overline{U}^{-0.16}$, $\sigma \sim \overline{U}^{-0.04}$



Figure 8: Dependence of reduced velocity \overline{V} and selection parameter σ on flow Reynolds number



Figure 9: Dependence of reduced tip radius $\overline{\rho}$ on reduced flow velocity \overline{U}

Growth of the dendrite tip in the parallel flow

Sidebranching in the parallel flow



Figure 10: Growth of side branches, $\Delta = 0.7, 15\varepsilon = 0.15$. Reduced flow velocity $\overline{U} = 0.01$ (a) and $\overline{U} = 0.04$ (b)

At large flow velocities, oscillations of tip velocity were observed, accompanied by the enhanced growth of side branches.

Conclusions

The main results are

- Simulations showed strong influence of external shear flow on the seaweed growth (Fig. 2)
- Influence of parallel flow on the operation state of dendrite tip was investigated quantitatively (Fig. 4–9)
- Simulations demonstrated onset of velocity oscillations and enhancement of sidebranching under parallel flow (Fig. 10)